

A Simplified Treatment of Ambipolar Diffusion?



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- Ambipolar diffusion (AD) is the process of ions and neutrals drifting apart in a partially ionized plasma (PIP),
 - Allowing the magnetic flux to diffuse out of compressed, low-ionization regions.
- AD is generally thought to be important for the collapse of magnetically supported molecular cloud cores.
- Its numerical calculation often involves very short timesteps.
 - Making numerical simulations expensive.
 - Often solved by advancing the induction equation alone on its own timestep (e.g., Mac Low+95)
- Is there a way to circumvent this technical difficulty?

AD basics:

- The relevant equations:

- Momentum equation:

$$\frac{\partial \mathbf{u}_n}{\partial t} + \mathbf{u}_n \cdot \nabla \mathbf{u}_n = -\nabla P_n + (\nabla \times \mathbf{B}) \times \mathbf{B}$$

Lorentz force

- Induction equation:

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{u}_n \times \mathbf{B}) + \nabla \times \left\{ \frac{\mathbf{B}}{4\pi\gamma\rho_n\rho_i} \times [\mathbf{B} \times (\nabla \times \mathbf{B})] \right\}$$

Flux freezing

AD

with ρ_n = mass density of neutrals; ρ_i = mass density of ions
 u_n = neutrals bulk velocity.

- AD has two simple limiting behaviors:

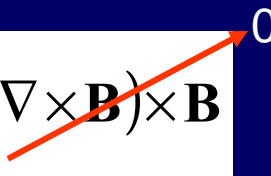
- Flux-freezing (ideal MHD) when AD is negligible.

$$\frac{\partial \mathbf{u}_n}{\partial t} + \mathbf{u}_n \cdot \nabla \mathbf{u}_n = -\nabla P_n + (\nabla \times \mathbf{B}) \times \mathbf{B}$$

$$\frac{\partial \mathbf{B}}{\partial t} \approx \nabla \times (\mathbf{u}_n \times \mathbf{B})$$

- Decoupled HD + uniform B when AD is strong.

$$\frac{\partial \mathbf{B}}{\partial t} \approx \nabla \times \left\{ \frac{\mathbf{B}}{4\pi\gamma\rho_n\rho_i} \times [\mathbf{B} \times (\nabla \times \mathbf{B})] \right\} \Rightarrow \mathbf{B} \rightarrow \text{cst.}$$

$$\frac{\partial \mathbf{u}_n}{\partial t} + \mathbf{u}_n \cdot \nabla \mathbf{u}_n = -\nabla P_n + (\nabla \times \mathbf{B}) \times \mathbf{B}$$


- *In the latter case, one is integrating frantically* (timescale for AD is \ll than flux-freezing term) *just to decouple the field!*

- Is it possible to take advantage of these limiting behaviors to avoid direct integration with small timesteps?
- A case where it works: cooling (VS et al. 2007, ApJ, 657, 870)
- The internal energy equation per unit mass including cooling ($n\Lambda$) and heating (Γ) specific rates is:

$$\frac{\partial e}{\partial t} + \mathbf{u} \cdot \nabla e = -(\gamma - 1)e \nabla \cdot \mathbf{u} + \Gamma - n\Lambda$$

- When heating and cooling dominate, this becomes

$$\frac{de}{dt} \approx \Gamma - n\Lambda \quad \Rightarrow \quad n\Lambda \rightarrow \Gamma \quad (\text{approach to thermal equilibrium})$$

- The thermal equilibrium condition

$$n\Lambda(n, e) = \Gamma(n, e)$$

together with the equation of state

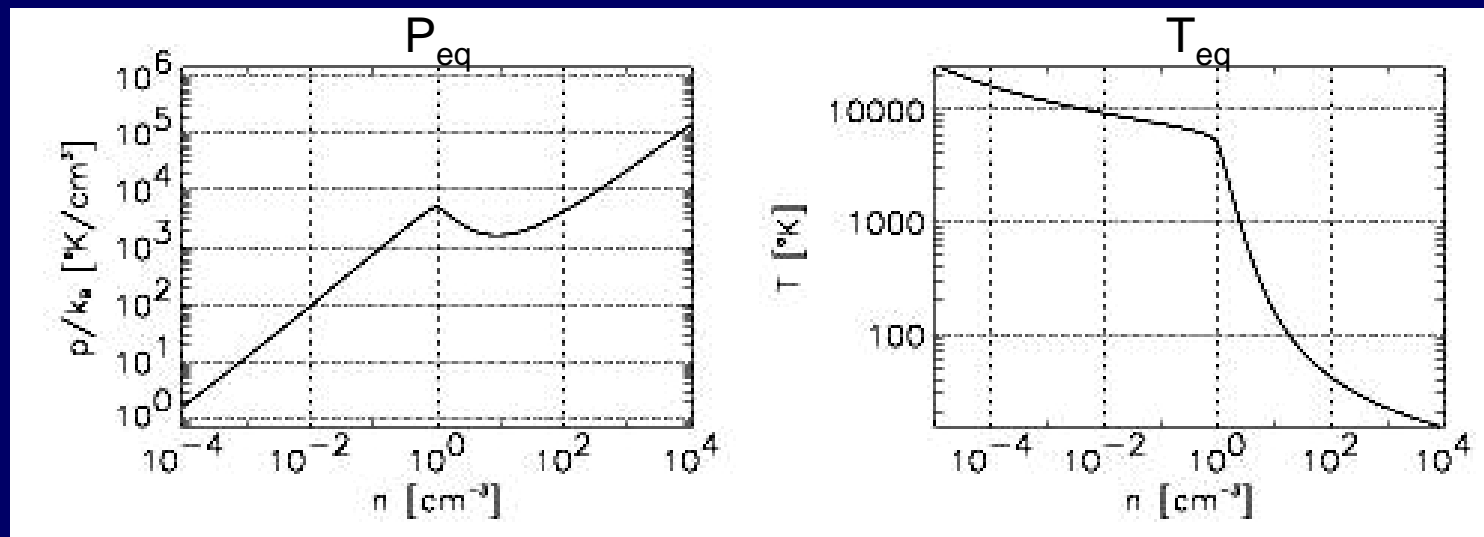
$$P = (\gamma - 1)\rho e,$$

allow one to eliminate P to get

$$e_{\text{eq}} = e_{\text{eq}}(\rho).$$

Thermal equilibrium; ($e \propto T_{\text{eq}}$).

- Can use a (perhaps piecewise) polytropic approximation:



- The timescale to reach equilibrium is

$$\tau = \frac{e - e_{\text{eq}}}{n\Lambda - \Gamma}$$

Cooling time

- Can be *much* shorter than dynamical timescale, especially at high n .
- But only effect is to approach thermal equilibrium.
 - A polytropic approximation would suffice.

- Then can “pre-integrate” over HD step Δt :

$$(e - e_{\text{eq}})^{-1} d(e - e_{\text{eq}}) = -\frac{n\Lambda - \Gamma}{e - e_{\text{eq}}} dt = -\frac{dt}{\tau}$$

\Rightarrow

$$e(t + \Delta t) = e_{\text{eq}} + [e(t) - e_{\text{eq}}] \exp\left(-\frac{\Delta t}{\tau}\right).$$

- This has the correct limiting behavior:

$$\text{If } \Delta t \gg \tau, \quad e(t + \Delta t) \rightarrow e_{\text{eq}};$$

$$\text{if } \Delta t \ll \tau, \quad e(t + \Delta t) \rightarrow e(t) - \Delta t(n\Lambda - \Gamma)$$

- Is a similar procedure possible for the AD?
 - How to interpolate between ideal MHD and HD + uniform B?
 - Can introduce weight factors into the Lorentz force and AD terms?

Thank you!