

# Dynamics of rain blobs: fully and partially ionised plasmas

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## Some observational facts

- 1 Acceleration typically much **smaller than that of gravity** (projected along the blob trajectory).
- 2 Blobs often reach a **constant velocity**, in the range 40–120 km s<sup>-1</sup>.
- 3 De Groof et al. (2005), using H $\alpha$  and He 304 Å observations: blob made of neutral and ionised fractions that fall together. **Strong (dynamical?) coupling between charged particles and neutrals.**

We consider a simple configuration and perform hydrodynamic numerical simulations to investigate the kinematics and dynamics of coronal rain blobs.



# Model: assumptions (FIP)

## We include:

- Fully ionised plasma (FIP) made of H,
- hydrodynamic time evolution of a formed blob in the loop environment,
- vertical motions,
- adiabatic changes of state.

## We do not include:

- The neutral fraction of the plasma,
- the blob formation process,
- path curvature,
- radiation, conduction, cooling,
- transition region and chromosphere,
- magnetic field.

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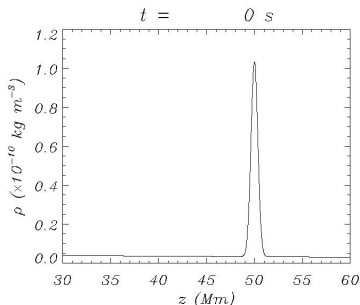


# Model: initial state

- Background: isothermal loop, vertically stratified by gravity:

$$p(z, t = 0) = p_0 e^{-z/H}, \quad \rho(z, t = 0) = \rho_0 e^{-z/H}, \quad H = \frac{2RT_0}{g}$$

- $z$ -axis is vertical;  $T_0$  is the loop temperature;  $R = k_B/m_p$  is the gas constant.
- The blob is represented by a localised **density enhancement**.
- In the initial state the vertical speed is zero.



Parameter values:

$$H = 120 \text{ Mm}$$

$$\rho_0 = 5 \times 10^{-12} \text{ kg m}^{-3}$$

$$\rho_{b0} = 10^{-10} \text{ kg m}^{-3}$$

$$z_0 = 50 \text{ Mm}$$

# Model: equations

- Three unknowns  $\rho$ ,  $p$ ,  $v$ .
- The unknowns only depend on  $t$  and  $z$ .

- Mass continuity:

$$\frac{\partial \rho}{\partial t} = -v \frac{\partial \rho}{\partial z} - \rho \frac{\partial v}{\partial z}$$

- $z$ -component of the momentum equation:

$$\rho \frac{\partial v}{\partial t} = -\rho v \frac{\partial v}{\partial z} - \frac{\partial p}{\partial z} - g\rho$$

- Energy equation:

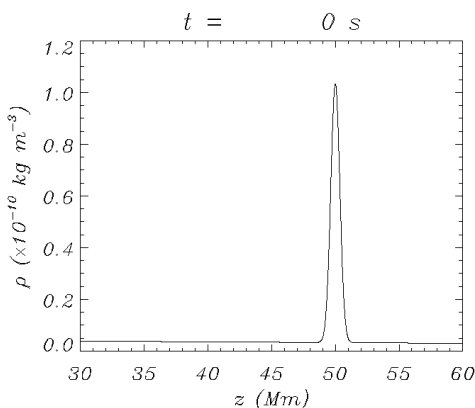
$$\frac{\partial p}{\partial t} = -v \frac{\partial p}{\partial z} - \gamma p \frac{\partial v}{\partial z}$$

# Model: numerical method

- The time-dependent PDEs are solved using the PDE2D code with Galerkin's method.
- We are interested in the height range  $0 \leq z \leq 60$  Mm.
- Open boundaries at the top and bottom of the system.
- No artificial viscosity/dissipation added.
- Grid spacing  $\simeq 10$  km.
- Simulation time: 1000 s.

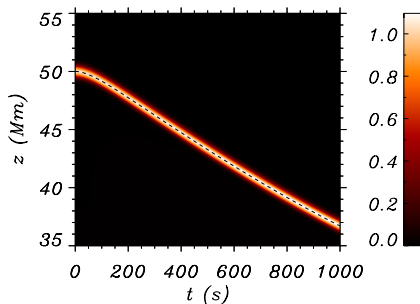
# Density

- The condensation keeps its density and shape as it falls.
- If the blob were to fall with the acceleration of gravity, it would reach  $z = 0$  at  $t \simeq 540$  s.



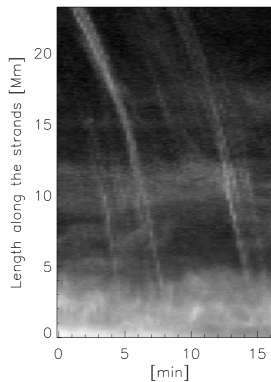
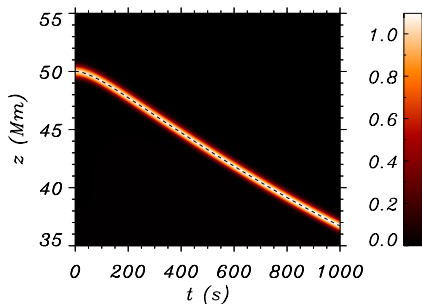
# Blob height

- Height–time diagram of density (colour scale in units of  $10^{-10} \text{ kg m}^{-3}$ ).
- Initial acceleration phase followed by constant speed phase.



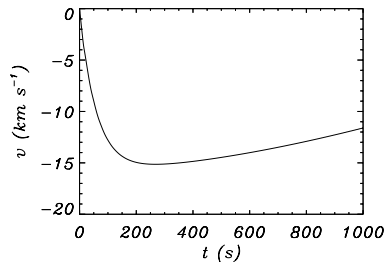
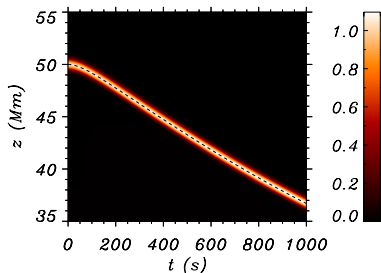
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Antolin et al. (2010)

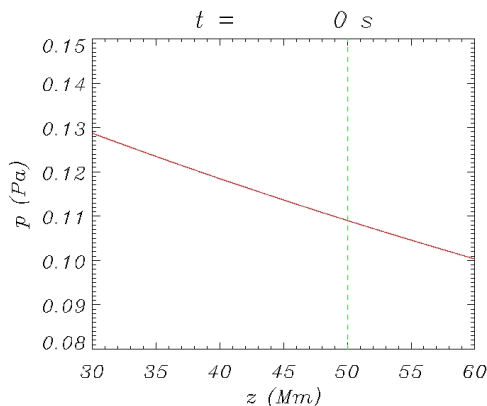
# Height and blob velocity



- Initial acceleration phase lasts  $\simeq 200$   $s$ .
- Blob then attains a roughly constant speed  $\simeq 15$   $km\ s^{-1}$ .

# Pressure

- Sound wave generated at the initial blob position (sound speed  $\simeq 235$  km s $^{-1}$ ).
- A strong pressure gradient is generated at the blob position.
- Established pressure gradient travels with the blob (its position is given by the green line).

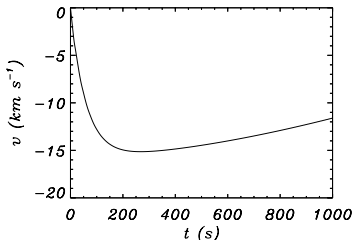
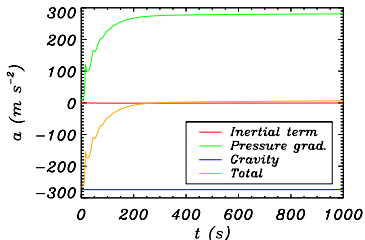


- Pressure structure under and above the blob is rearranged.



# Blob acceleration

$$\rho \frac{\partial v}{\partial t} = -\rho v \frac{\partial v}{\partial z} - \frac{\partial p}{\partial z} - g\rho \quad \Rightarrow \quad \frac{\partial v}{\partial t} = -v \frac{\partial v}{\partial z} - \frac{1}{\rho} \frac{\partial p}{\partial z} - g$$



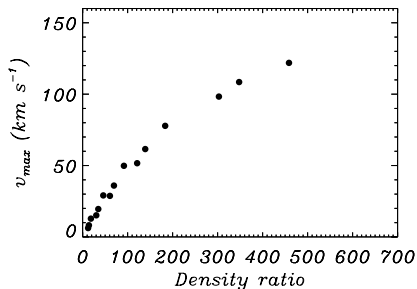
- Inertial term is negligible.
- Pressure gradient builds up and counteracts gravity at  $t \simeq 200$  s.

# Maximum descending speed

Similar dynamics are obtained for other parameter values.

- The **density ratio** (initial blob density to local loop density) determines the maximum falling speed.
- The local loop density depends on
  - initial blob height ( $z_0$ ),
  - density at base of loop ( $\rho_0$ ),
  - the vertical scale height ( $H \propto T_0$ ).

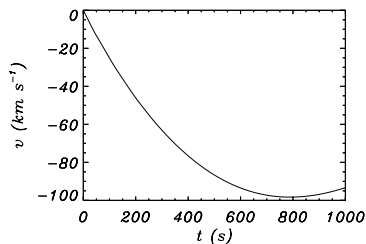
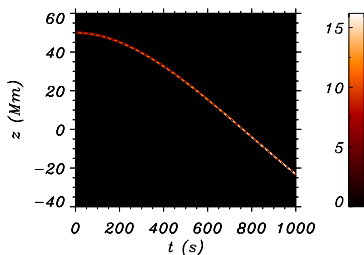
$$\rho(z, t = 0) = \rho_0 e^{-z/H}$$



- Observed blob velocities require blob to loop density ratios as large as 500 or more.

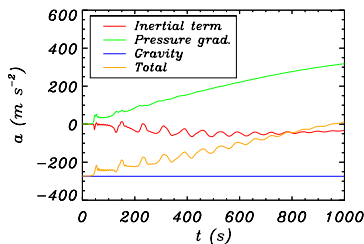
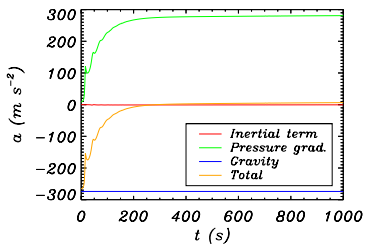
# Larger blob density: height and velocity

- Blob density is 10 times larger.
- Height–time diagram is much more parabolic.

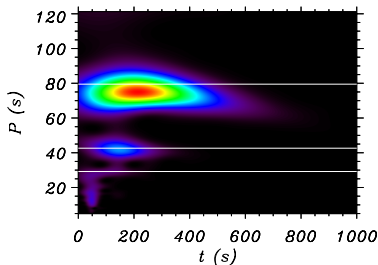


- Initial acceleration phase lasts much longer ( $\simeq 800 s$  compared to  $\simeq 200 s$ ).
- Maximum falling velocity is much higher ( $\simeq 100 km s^{-1}$  compared to  $\simeq 15 km s^{-1}$ )

# Larger blob density: acceleration

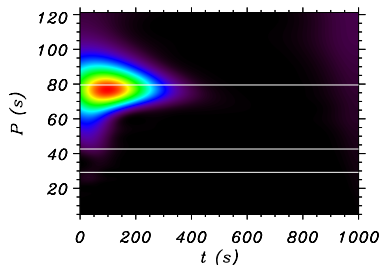


- Oscillations in the inertial term acceleration:  $v \frac{\partial v}{\partial z}$ .
- Wavelet diagram of  $\frac{\partial v}{\partial z}$ .
- Periods:  $\simeq 75$  s,  $\simeq 42$  s,  $\simeq 28$  s.

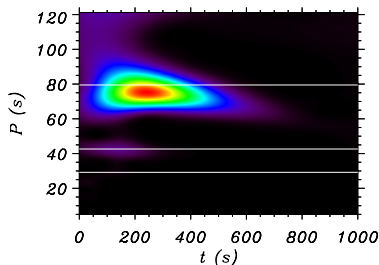


# Larger blob density: sound wave emission

- Blob density and velocity also display damped oscillations.
- These oscillations in  $\rho$  and  $v$  are revealed after fitting and subtracting a third degree polynomial + wavelet analysis.
- Horizontal lines: theoretical periods of leaky sound waves of static blob with  $g = 0$ .



Blob velocity



Blob density

# Model: assumptions (PIP)

## We include:

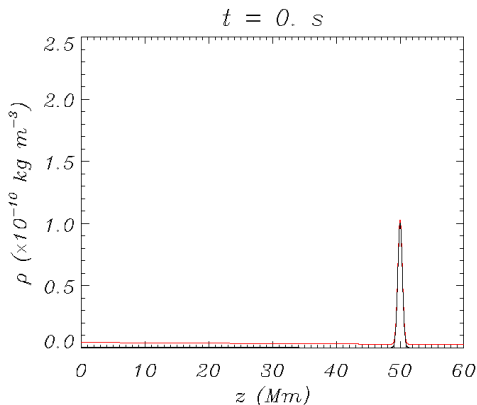
- The blob is partially ionised.
- The loop also contains a small fraction of neutrals (to avoid numerical problems).
- Hydrodynamic equations for the time evolution of two fluids (charged particles and neutrals) **coupled by friction**.

## We do not include:

- Ionisation and recombination.

# Model: initial state

- Red = charged particles, black = neutrals.



# Model: initial state

- Background: isothermal loop, vertically stratified by gravity:

$$p_{ie}(z, t = 0) = p_{ie0} e^{-z/H_i}, \quad \rho_i(z, t = 0) = \rho_{i0} e^{-z/H_i}, \quad H_i = \frac{2RT_0}{g}$$

$$p_n(z, t = 0) = p_{n0} e^{-z/H_n}, \quad \rho_n(z, t = 0) = \rho_{n0} e^{-z/H_n}, \quad H_n = \frac{RT_0}{g}$$

- Parameter values:

$$H_i = 120 \text{ Mm}$$

$$\rho_{i0} = 4.5 \times 10^{-12} \text{ kg m}^{-3}$$

$$\rho_{bi0} = 10^{-10} \text{ kg m}^{-3}$$

$$z_0 = 50 \text{ Mm}$$

$$H_n = 60 \text{ Mm}$$

$$\rho_{n0} = 0.5 \times 10^{-12} \text{ kg m}^{-3}$$

$$\rho_{bn0} = 10^{-10} \text{ kg m}^{-3}$$



# Model: equations

- Six unknowns  $\rho_i$ ,  $p_{ie}$ ,  $v_i$  and  $\rho_n$ ,  $p_n$ ,  $v_n$ .

- Mass continuity equations:

$$\frac{\partial \rho_i}{\partial t} = -v_i \frac{\partial \rho_i}{\partial z} - \rho_i \frac{\partial v_i}{\partial z} \qquad \frac{\partial \rho_n}{\partial t} = -v_n \frac{\partial \rho_n}{\partial z} - \rho_n \frac{\partial v_n}{\partial z}$$

- $z$ -component of the momentum equations:

$$\begin{aligned} \rho_i \frac{\partial v_i}{\partial t} &= -\rho_i v_i \frac{\partial v_i}{\partial z} - \frac{\partial p_{ie}}{\partial z} - g \rho_i & \rho_n \frac{\partial v_n}{\partial t} &= -\rho_n v_n \frac{\partial v_n}{\partial z} - \frac{\partial p_n}{\partial z} - g \rho_n \\ &- \alpha_{in}(v_i - v_n) & &+ \alpha_{in}(v_i - v_n) \end{aligned}$$

- Energy equations:

$$\begin{aligned} \frac{\partial p_{ie}}{\partial t} &= -v_i \frac{\partial p_{ie}}{\partial z} - \gamma p_{ie} \frac{\partial v_i}{\partial z} & \frac{\partial p_n}{\partial t} &= -v_n \frac{\partial p_n}{\partial z} - \gamma p_n \frac{\partial v_n}{\partial z} \\ &+ (\gamma - 1) Q_i^{in} & &- (\gamma - 1) Q_i^{in} \end{aligned}$$

# Model: equations

- $Q_i^{in}$  = heating of ions due to collisions with neutrals.

$$Q_i^{in} = \alpha_{in} \left[ \frac{1}{2} (v_n - v_i)^2 + 3R(T_n - T_i) \right]$$

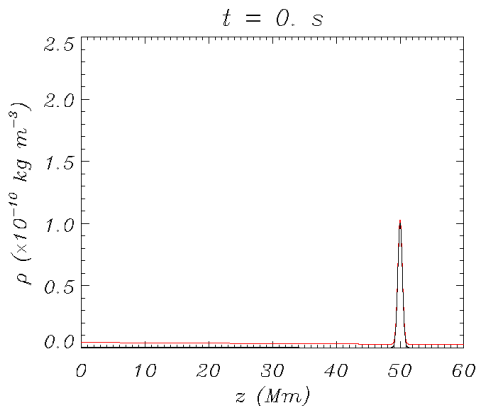
- The temperatures of ions and neutrals are obtained from the perfect gas law

$$p_{ie} = 2\rho_i R T_i \quad p_n = \rho_n R T_n$$

- The friction coefficient  $\alpha_{in}$  depends on the density and temperature of charged particles and neutrals.  
It must be **computed at each position at each time step.**

# Density

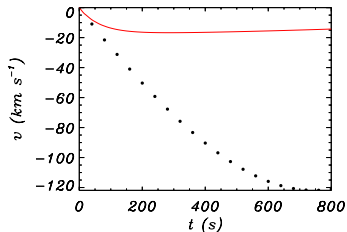
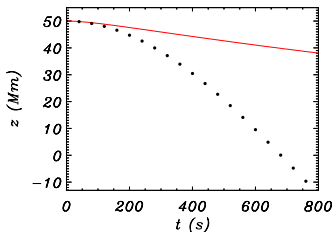
- No friction force ( $\alpha_{in} = 0$ ).
- The neutral fraction of the blob (black line) falls faster than the **ionised fraction** because of its larger density ratio.



# Height vs. time and velocity

Black = neutrals; red = ions.

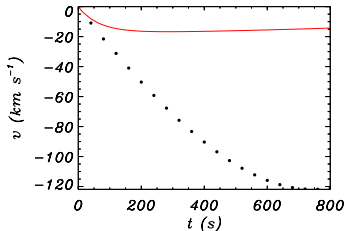
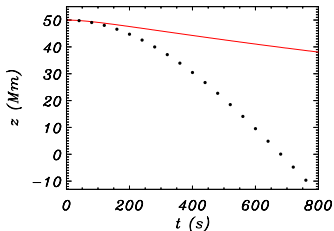
$$\alpha_{in} = 0$$



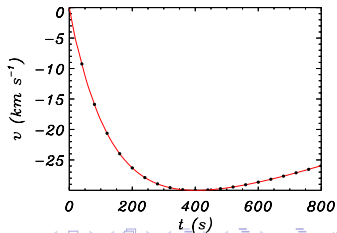
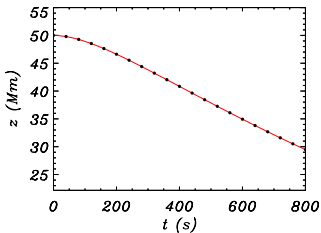
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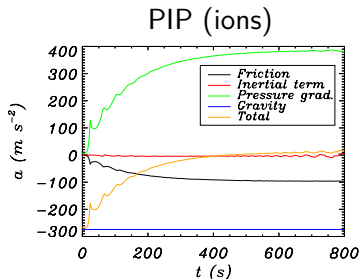
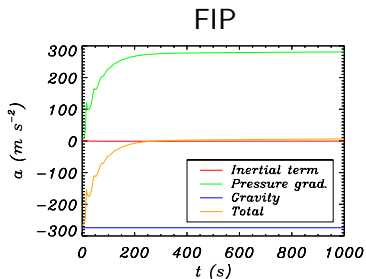
$$\alpha_{in} = 0$$



$$\alpha_{in} \neq 0$$



# Acceleration

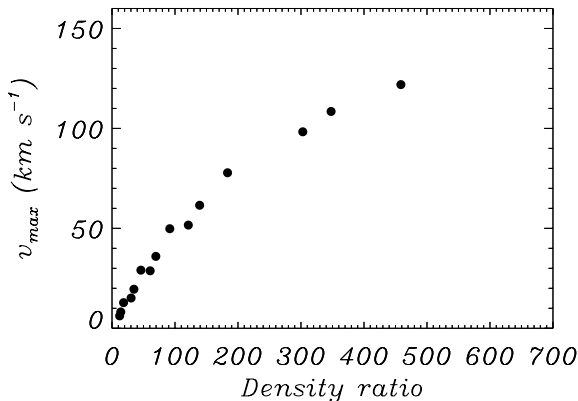


- Friction with neutrals gives rise to downward force that accelerates ions.

# Maximum descending speed

Black circles = fully ionised plasma.

Blue squares = partially ionised plasma.

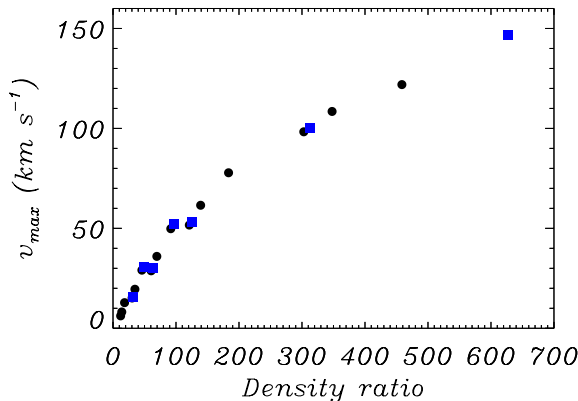


- The **density ratio** still determines the maximum falling speed.

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- The **density ratio** still determines the maximum falling speed.



# Conclusions

- Strong dynamical coupling of charged particles and neutrals. Main results do not change when going from FIP to PIP.
- Presence of mass condensation leads to a **pressure gradient that opposes gravity**.
- “Small” density ratio: acceleration phase followed by more or less **constant speed**.
- “Large” density ratio: acceleration phase lasts much longer.
- **Maximum falling speed increases with density ratio**. Very good correlation between the two quantities.
- Damped oscillations because of sound wave emission. Detectable as **periodic changes of the blob density and velocity?**